

# An investigation of structural parameters and magnetic and optical properties of $\text{EuLn}_2\text{Q}_4$ ( $\text{Ln} = \text{Tb–Lu}$ , $\text{Q} = \text{S, Se}$ )

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## Abstract

$\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) has been synthesized using  $\text{Sb}_2\text{Q}_3$  ( $\text{Q} = \text{S, Se}$ ) fluxes at 1000 °C. These compounds crystallize in a  $\text{CaFe}_2\text{O}_4$ -type three-dimensional channel structure that is built from edge-shared double rutile chains of  $[\text{LnQ}_6]$  octahedra running down the *b*-axis. Each double chain is connected at the vertices to four other double chains to form open channels where bicapped trigonal prismatic  $\text{Eu}^{2+}$  ions reside. All of these compounds show antiferromagnetic ordering with Neel temperatures,  $T_N \sim 3–4$  K. The optical band gaps for  $\text{EuTb}_2\text{Se}_4$ ,  $\text{EuDy}_2\text{Se}_4$ ,  $\text{EuHo}_2\text{Se}_4$ ,  $\text{EuEr}_2\text{Se}_4$ ,  $\text{EuTm}_2\text{Se}_4$ ,  $\text{EuYb}_2\text{Se}_4$ ,  $\text{EuLu}_2\text{Se}_4$ , and  $\text{EuYb}_2\text{S}_4$  are found to be 2.0, 1.8, 1.8, 1.7, 1.8, 1.3, 1.7, and 1.6 eV, respectively.

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## 1. Introduction

A large number of ternary compounds with a formula  $\text{ALn}_2\text{Q}_4$ , where  $A = \text{Sr, Ba, Eu}$ ;  $\text{Ln} = \text{La–Lu, Y, Sc}$ ;  $\text{Q} = \text{O, S, Se, Te}$ , crystallize in the  $\text{CaFe}_2\text{O}_4$ -type [1] structure [2–11]. These systems have been extensively studied owing to their interesting optical and magnetic properties. Some of these compounds are potential IR window materials [7].  $\text{SrLn}_2\text{O}_4$  ( $\text{Ln} = \text{Dy, Ho, Er, Tm, Yb}$ ) [10] and  $\text{BaLn}_2\text{O}_4$  ( $\text{Ln} = \text{Pr, Nd, Sm–Ho}$ ) [11] with a triangle-based array of magnetic *Ln* ions were recently recognized as potentially geometrically frustrated systems. Compared to  $\text{SrLn}_2\text{Q}_4$  or  $\text{BaLn}_2\text{Q}_4$ ,  $\text{Eu}^{2+}$ -based phases are more magnetically interesting.  $\text{Eu}^{2+}$  cations have fairly large magnetic moments ( $7.94 \mu_B$ ) and potentially strong magnetic couplings between each other, which usually leads to long-range magnetic ordering [12,13]. Many  $\text{EuLn}_2\text{Q}_4$  ( $\text{Q} = \text{S, Se}$ ) compounds have been characterized by powder X-ray diffraction and magnetic measurements [3,4]. They have shown antiferromagnetic

transitions, metamagnetic transitions, and ferromagnetic transitions, depending on the choice of  $\text{Ln}^{3+}$ . To our knowledge,  $\text{EuEr}_2\text{S}_4$  is the only compound of the  $\text{EuLn}_2\text{Q}_4$  ( $\text{Q} = \text{S, Se}$ ) [5] series examined with single X-ray diffraction. Selenide analogs are less studied compared to  $\text{EuLn}_2\text{S}_4$ . Here we present the syntheses, single crystal structures, optical, and magnetic properties of  $\text{EuLn}_2\text{Se}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ).  $\text{EuYb}_2\text{S}_4$  is also included in the discussion for comparison.

## 2. Experimental

### 2.1. Starting materials

Eu (99.9%, Alfa-Aesar), Tb (99.9%, Alfa-Aesar), Dy (99.9%, Alfa-Aesar), Ho (99.9%, Alfa-Aesar), Er (99.9%, Alfa-Aesar), Tm (99.9%, Alfa-Aesar), Yb (99.9%, Alfa-Aesar), Lu (99.9%, Alfa-Aesar), Se (99.5%, Alfa-Aesar), S (99.5%, Alfa-Aesar), and Sb (99.5%, Alfa-Aesar) were used as received. The  $\text{Sb}_2\text{Q}_3$  ( $\text{Q} = \text{S, Se}$ ) fluxes were prepared from the direct reaction of the elements in sealed fused-silica ampoules at 850 °C.

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## 2.2. Syntheses

$\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) were obtained from the reaction of  $\text{Eu}$ ,  $\text{Ln}$ ,  $\text{Q}$ , and  $\text{Sb}_2\text{Q}_3$  ( $\text{Q} = \text{S, Se}$ ) in a molar ratio of 1:2:4:0.5. All of the mixtures were loaded into fused-silica ampoules in an Ar-filled glovebox. The ampoules were sealed under vacuum and heated in a programmable tube furnace using the following profile: 2 °C/min to 500 °C (held for 1 h), 0.5 °C/min to 1000 °C (held for 7 d), 0.05 °C/min to 550 °C (held for 2 d), and 0.5 °C/min to 24 °C. The products included high yields of red or black crystals of  $\text{EuLn}_2\text{Q}_4$  and unreacted  $\text{Sb}_2\text{Q}_3$  blocks. For  $\text{EuLn}_2\text{Se}_4$  ( $\text{Ln} = \text{Tb, Dy, Ho}$ ), the yield of desired products becomes low. Instead, there were large amounts of  $\text{EuSbSe}_3$  [13] and some unidentified black microcrystalline ‘ $\text{EuLnSbSe}$ ’ phases. EDX experiments have shown that the crystals of ‘ $\text{EuLnSbSe}$ ’ include all four elements, but they are too small for single-crystal X-ray diffraction experiments.  $\text{EuLn}_2\text{Q}_4$  crystals were isolated manually for physical property measurements. All of the compounds are stable in air. Powder X-ray diffraction measurements were used to confirm phase purity by comparing the powder patterns calculated from the single crystal X-ray structures with the experimental data. Semi-quantitative SEM/EDX analyses were performed using JEOL 840/Link Isis or JEOL JSM-7000F instruments.  $\text{Eu}$ ,  $\text{Ln}$ , and  $\text{Q}$  percentages were calibrated against standards.  $\text{Sb}$  was not detected in the crystals. The  $\text{Eu}:\text{Ln}:\text{Q}$  ratios were determined to be approximately 1:2:4 from EDX analyses.

## 2.3. Crystallographic studies

Single crystals of  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) were mounted on glass fibers with epoxy and optically

aligned on a Bruker APEX single-crystal X-ray diffractometer using a digital camera. Initial intensity measurements were performed using graphite monochromated  $\text{MoK}\alpha$  ( $\lambda = 0.71073 \text{ \AA}$ ) radiation from a sealed tube and a monocapillary collimator. SMART (v 5.624) was used for preliminary determination of the cell constants and data collection control. The intensities of reflections of a sphere were collected by a combination of three sets of exposures (frames). Each set had a different  $\phi$  angle for the crystal and each exposure covered a range of 0.3° in  $\omega$ . A total of 1800 frames were collected with exposure times per frame of 10 or 20 s depending on the crystal.

For  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ), determinations of integrated intensities and global refinement were performed with the Bruker SAINT (v 6.02) software package using a narrow-frame integration algorithm. These data were treated first with a face-index numerical absorption correction using XPREP [14], followed by a semi-empirical absorption correction using SADABS [15]. The program suite SHELXTL (v 6.12) was used for space group determination (XPREP), direct methods structure solution (XS), and least-squares refinement (XL) [14]. The final refinements included anisotropic displacement parameters for all atoms and secondary extinction. Some crystallographic details are given in Table 1. As an example, atomic coordinates and equivalent isotropic displacement parameters for  $\text{EuYb}_2\text{S}_4$  are given in Table 2. Additional crystallographic details can be found in the Supporting Information. Further details of the crystal structure investigation may be obtained from the Fachinformativzentrum Karlsruhe, D-76344 Eggenstein-Leopoldshafen, Germany (fax: +49 7247 808 666; Email: [crysdata@fiz-karlsruhe.de](mailto:crysdata@fiz-karlsruhe.de)) on quoting the depository numbers CSD 418588, 418589, 418590, 418591, 418592, 418593, 418594, and 418595.

Table 1  
Crystallographic data for  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ )

Formula	$\text{EuTb}_2\text{Se}_4$	$\text{EuDy}_2\text{Se}_4$	$\text{EuHo}_2\text{Se}_4$	$\text{EuEr}_2\text{Se}_4$	$\text{EuTm}_2\text{Se}_4$	$\text{EuYb}_2\text{Se}_4$	$\text{EuLu}_2\text{Se}_4$	$\text{EuYb}_2\text{S}_4$
$F_w$	785.64	792.80	797.66	802.32	805.66	813.88	817.74	626.28
Color	Red	Red	Red	Red	Red	Black	Red	Dark red
Crystal system	Orthorhombic	Orthorhombic	Orthorhombic	Orthorhombic	Orthorhombic	Orthorhombic	Orthorhombic	Orthorhombic
Space group	$Pnma$ (no. 62)	$Pnma$ (no. 62)	$Pnma$ (no. 62)	$Pnma$ (no. 62)	$Pnma$ (no. 62)	$Pnma$ (no. 62)	$Pnma$ (no. 62)	$Pnma$ (no. 62)
$a$ (Å)	12.429(2)	12.442(1)	12.4258(7)	12.4066(7)	12.3874(8)	12.3591(9)	12.3489(9)	11.809(1)
$b$ (Å)	4.1413(5)	4.1224(3)	4.1027(2)	4.0864(2)	4.0737(3)	4.0612(3)	4.0470(3)	3.9194(3)
$c$ (Å)	14.915(2)	14.866(1)	14.8096(8)	14.7596(8)	14.7119(9)	14.664(1)	14.6174(1)	14.086(1)
$V$ (Å <sup>3</sup> )	767.7(3)	762.5(2)	755.0(1)	748.3(1)	742.4(1)	736.0(2)	730.5(2)	651.9(2)
$Z$	4	4	4	4	4	4	4	4
$T$ (K)	193	193	193	193	193	193	193	193
$\lambda$ (Å)	0.71073	0.71073	0.71073	0.71073	0.71073	0.71073	0.71073	0.71073
$\rho_{\text{calcd}}$ (g cm <sup>-3</sup> )	6.797	6.906	7.018	7.122	7.208	7.345	7.435	6.381
$\mu$ (cm <sup>-1</sup> )	451.43	465.00	481.28	498.42	515.30	532.82	551.06	391.00
$R$ ( $F$ ) <sup>a</sup>	0.0304	0.0298	0.0219	0.0214	0.0248	0.0211	0.0190	0.0208
$R_w$ ( $F_o^2$ ) <sup>b</sup>	0.0656	0.0667	0.0490	0.0501	0.0606	0.0512	0.0448	0.0471

$$^a R(F) = \frac{\sum |F_o| - |F_c|}{\sum |F_o|}$$

$$^b R_w(F_o^2) = \frac{\sum [w(F_o^2 - F_c^2)^2]}{\sum wF_o^4}^{1/2}$$

Table 2  
Atomic coordinates and equivalent isotropic displacement parameters for  $\text{EuYb}_2\text{S}_4$

Atom (site)	x	y	z	$U_{\text{eq}} (\text{\AA}^2)^a$
Eu	0.75797(4)	$\frac{1}{4}$	0.66230(3)	0.0083(1)
Yb1	0.41916(3)	$\frac{1}{4}$	0.10025(2)	0.0078(1)
Yb2	0.43311(3)	$\frac{1}{4}$	0.61031(3)	0.0088(1)
S1	0.2080(2)	$\frac{1}{4}$	0.1735(2)	0.0086(4)
S2	0.1293(2)	$\frac{1}{4}$	0.4729(1)	0.0071(4)
S3	0.5226(2)	$\frac{1}{4}$	0.7841(2)	0.0094(4)
S4	0.4131(2)	$\frac{1}{4}$	0.4219(2)	0.0075(4)

<sup>a</sup> $U_{\text{eq}}$  is defined as one-third of the trace of the orthogonalized  $U_{ij}$  tensor.

#### 2.4. Powder X-ray diffraction

Powder X-ray diffraction patterns were collected with a Rigaku Miniflex powder X-ray diffractometer using  $\text{CuK}\alpha$  ( $\lambda = 1.54056 \text{ \AA}$ ) radiation.

#### 2.5. Magnetic susceptibility measurements

Magnetism data were measured on powders obtained from ground crystals in gelcap sample holders with a Quantum Design MPMS 7T magnetometer/susceptometer between 2 and 300 K and in applied fields up to 7 T. DC susceptibility measurements were made under zero-field-cooled conditions with an applied field of 0.1 T. Susceptibility values were corrected for the sample diamagnetic contribution according to Pascal's constants [16] as well as for the sample holder diamagnetism.  $\theta_p$  values were obtained from extrapolation fits between 100 and 300 K.

#### 2.6. UV–vis–NIR diffuse reflectance spectroscopy

The diffuse reflectance spectra for  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) in the form of ground powders were measured from 200 to 2500 nm using a Shimadzu UV3100 spectrophotometer equipped with an integrating sphere attachment. The Kubelka–Munk function was used to convert diffuse reflectance data to absorption spectra [17].

### 3. Results and discussion

#### 3.1. Crystal structures

The  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) compounds adopt the  $\text{CaFe}_2\text{O}_4$ -type [1] three-dimensional channel structure, which has been detailed in other references [5–11]. These compounds crystallize in an orthorhombic system with a space group of  $Pnma$ . As shown in Fig. 1, the structure includes one crystallographically unique bicapped trigonal prismatic  $\text{Eu}^{2+}$  position and two octahedral  $\text{Ln}^{3+}$  sites.  $[\text{LnQ}_6]$  octahedra share edges with identical units

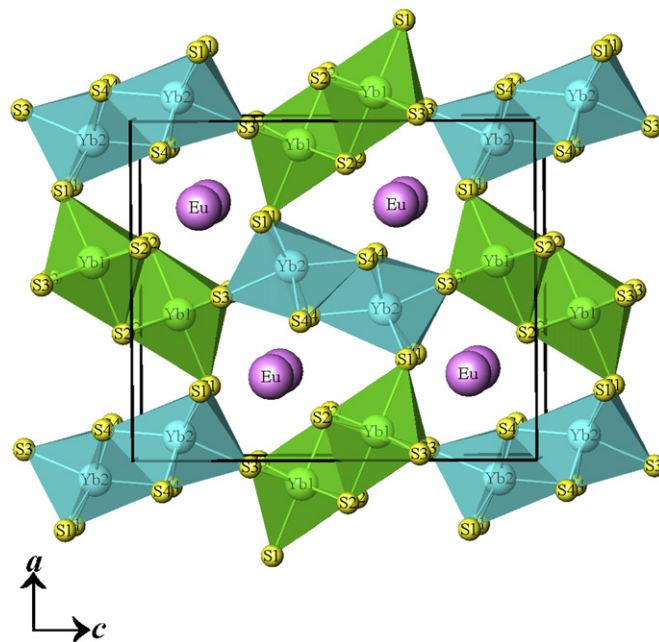


Fig. 1. A view of the three-dimensional channel structure of  $\text{EuYb}_2\text{S}_4$  down the  $b$ -axis.

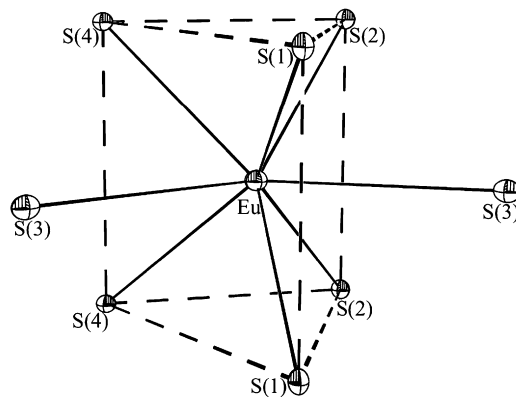


Fig. 2. Illustrations of the coordination environment for Eu ions in  $\text{EuYb}_2\text{S}_4$ .

both in the direction of chain propagation (the  $b$ -axis) and with adjacent chains to form double rutile chains. Each double chain is connected at the vertices to four other double chains to build open channels where  $\text{Eu}^{2+}$  ions reside. The coordination environment of  $\text{Eu}^{2+}$  ions is displayed in Fig. 2. Selected bond distances ( $\text{\AA}$ ) for  $\text{EuLn}_2\text{Q}_4$  are listed in Table 3.  $\text{Eu–Q}$  distances are in the range of 3.1469(7) and 3.438(2)  $\text{\AA}$  for selenides, and 3.039(2) and 3.266(2)  $\text{\AA}$  for the sulfide. They are comparable to Shannon's values for eight-coordinated  $\text{Eu}^{2+}$  ions [18]. Bond distances for  $[\text{LnSe}_6]$  octahedra are normal as well.

The unit cell volumes and  $\text{Eu–Eu}$  distances ( $b$ -axis) are plotted as a function of choice of  $\text{Ln}^{3+}$  in Fig. 3. These values decrease linearly with increasing atomic number because of the lanthanide contraction.

Table 3  
Selected bond distances (Å) for  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb-Lu}$ ;  $\text{Q} = \text{S, Se}$ )

Formula	$\text{EuTb}_2\text{Se}_4$	$\text{EuDy}_2\text{Se}_4$	$\text{EuHo}_2\text{Se}_4$	$\text{EuEr}_2\text{Se}_4$	$\text{EuTm}_2\text{Se}_4$	$\text{EuYb}_2\text{Se}_4$	$\text{EuLu}_2\text{Se}_4$	$\text{EuYb}_2\text{S}_4$
$\text{Eu-Q}(1) \times 2$	3.176(1)	3.176(1)	3.1747(7)	3.1740(8)	3.1754(9)	3.1701(8)	3.1730(7)	3.058(2)
$\text{Eu-Q}(2) \times 2$	3.183(1)	3.177(1)	3.1670(7)	3.1616(7)	3.1571(9)	3.1506(7)	3.1469(7)	3.039(2)
$\text{Eu-Q}(3)$	3.380(2)	3.392(1)	3.3933(9)	3.388(1)	3.377(1)	3.367(1)	3.3501(9)	3.215(2)
$\text{Eu-Q}(3)$	3.438(2)	3.419(1)	3.4024(9)	3.391(1)	3.387(1)	3.379(1)	3.3823(9)	3.266(2)
$\text{Eu-Q}(4) \times 2$	3.204(1)	3.205(1)	3.1980(7)	3.1945(7)	3.1920(9)	3.1848(7)	3.1831(7)	3.054(2)
$\text{Ln}(1)-\text{Q}(1)$	2.870(1)	2.864(1)	2.8566(8)	2.8458(9)	2.838(1)	2.8256(9)	2.8196(8)	2.699(2)
$\text{Ln}(1)-\text{Q}(2) \times 2$	2.8955(9)	2.8794(9)	2.8643(6)	2.8500(7)	2.8368(8)	2.8267(6)	2.8132(6)	2.718(1)
$\text{Ln}(1)-\text{Q}(2)$	2.872(1)	2.858(1)	2.8413(8)	2.8280(9)	2.817(1)	2.8042(9)	2.7944(8)	2.687(2)
$\text{Ln}(1)-\text{Q}(3) \times 2$	2.8012(9)	2.7940(9)	2.7812(6)	2.7722(6)	2.7643(7)	2.7551(6)	2.7473(6)	2.640(1)
$\text{Ln}(2)-\text{Q}(1) \times 2$	2.884(1)	2.8753(9)	2.8636(6)	2.8531(7)	2.8428(8)	2.8356(6)	2.8242(6)	2.722(2)
$\text{Ln}(2)-\text{Q}(3)$	2.837(1)	2.819(1)	2.8082(8)	2.7959(9)	2.786(1)	2.7751(9)	2.7653(8)	2.666(2)
$\text{Ln}(2)-\text{Q}(4) \times 2$	2.878(1)	2.8680(8)	2.8532(6)	2.8410(6)	2.8304(7)	2.8229(6)	2.8112(6)	2.710(1)
$\text{Ln}(2)-\text{Q}(4)$	2.839(1)	2.826(1)	2.8163(8)	2.803(1)	2.788(1)	2.777(1)	2.7658(9)	2.665(2)

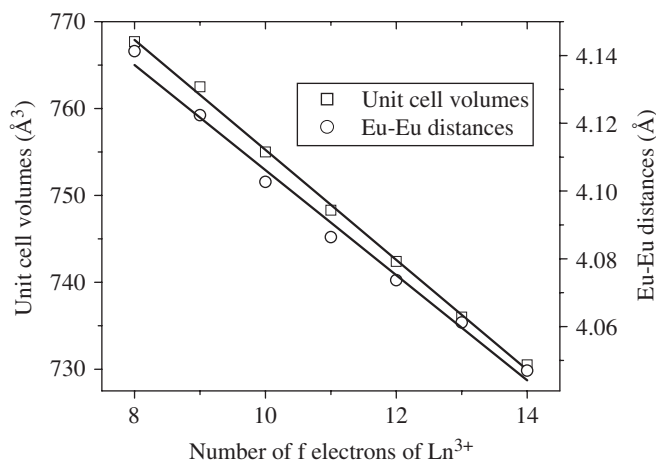


Fig. 3. Unit cell volumes ( $\text{Å}^3$ ) and Eu–Eu distances ( $\text{Å}$ ) vs. the number of f electrons of  $\text{Ln}^{3+}$  for  $\text{EuLn}_2\text{Se}_4$  ( $\text{Ln} = \text{Tb-Lu}$ ).

### 3.2. Magnetic properties

Magnetic susceptibilities of  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb-Lu}$ ;  $\text{Q} = \text{S, Se}$ ) have been measured on powder samples. Only the data for  $\text{EuYb}_2\text{Se}_4$  and  $\text{EuLu}_2\text{Se}_4$  are shown in Figs. 4 and 5. The magnetic behavior of other materials can be found in the Supporting Information (Figs. S1–S6). All of these compounds except  $\text{EuEr}_2\text{Se}_4$  and  $\text{EuHo}_2\text{Se}_4$  exhibit a similar transition below 5 K, which is relatively independent of the choice of  $\text{Ln}^{3+}$ . The transitions for  $\text{EuEr}_2\text{Se}_4$  and  $\text{EuHo}_2\text{Se}_4$  are anomalous, as both samples show a sharp antiferromagnetic transition on a large background of Curie-like Er or Ho paramagnetism. The transitions show no divergence in zero-field-cooled vs. field-cooled measurements, so they are ascribed to signify conventional antiferromagnetic ordering. As mentioned earlier, the triangular characteristics of the  $\text{Ln}^{3+}$  sublattice suppress the potential long-range ordering of the  $\text{Ln}^{3+}$  ions [10,11], even though this effect is expected to be less significant in sulfides and selenides than in oxides. In addition,  $T_N$  for  $\text{EuLu}_2\text{Se}_4$  is nearly identical to that of  $\text{EuYb}_2\text{Se}_4$  and  $\text{EuTm}_2\text{Se}_4$ , and  $\text{Lu}^{3+}$  ions carry no moment. However,

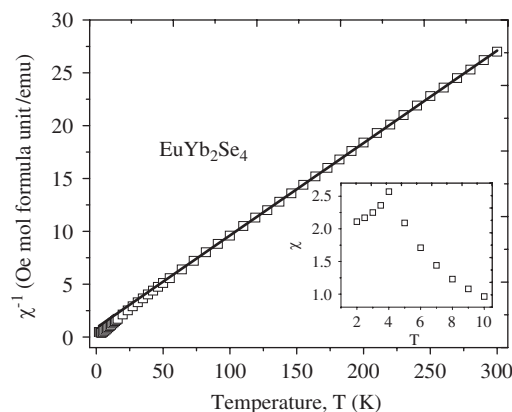


Fig. 4. Temperature dependence of the reciprocal molar magnetic susceptibility for  $\text{EuYb}_2\text{Se}_4$  under an applied magnetic field of 0.1 T between 2 and 300 K. The straight line represents the fit to the Curie–Weiss law in the range of 100–300 K. Inset shows the molar magnetic susceptibility at low temperature.

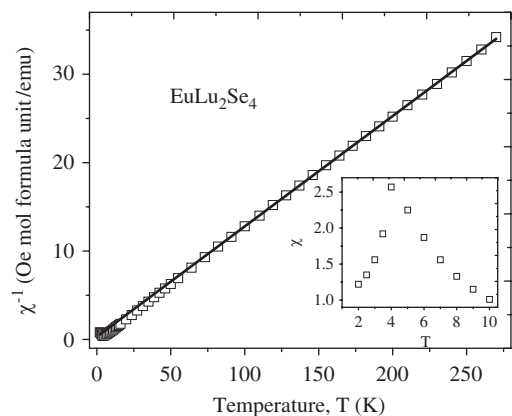


Fig. 5. Inverse molar magnetic susceptibility vs. temperature for  $\text{EuLu}_2\text{Se}_4$  under an applied magnetic field of 0.1 T between 2 and 300 K. The solid line represents the fit to the Curie–Weiss law in the range of 100–300 K. Inset shows the molar magnetic susceptibility at low temperature.

there is a weak inverse correlation between the  $\text{Eu}^{2+}$ -generated antiferromagnetic transition temperature and the total angular momentum quantum number of the  $\text{Ln}^{3+}$

ion. Clearly the observed magnetic ordering in these materials is due to the  $\text{Eu}^{2+}$  ions. The Eu–Eu coupling is brought about by Eu– $Q$ –Eu superexchange that is also exemplified by antiferromagnetic  $\text{Eu}_3\text{O}_4$  [19], which has a somewhat similar crystal structure.

The high-temperature susceptibilities of  $\text{EuLn}_2\text{Q}_4$  obey the Curie–Weiss law. Deviations can be found for most of these compounds at relatively low temperatures caused by a crystal-field splitting of the ground state for the  $\text{Ln}^{3+}$  cations. (Crystal field splitting effects will be negligible for  $\text{Eu}^{2+}$  ions, as the ground-state angular momentum quantum number for  $\text{Eu}^{2+}$  is  $J = S = 7/2$ , thus with no orbital contribution.) Table 4 lists the magnetic parameters for  $\text{EuLn}_2\text{Q}_4$ , which were obtained from fitting the data in the range of 100–300 K with the Curie–Weiss law. The experimental effective magnetic moments are all close to that expected for full moments of the free  $\text{Eu}^{2+}$  and  $\text{Ln}^{3+}$  ions [20]. Weiss temperatures ( $\theta_p$ ) are all negative, which is consistent with antiferromagnetic interactions between magnetic ions. Detailed magnetic properties including the hydrostatic pressure dependence of the antiferromagnetic ordering temperatures are discussed in another paper [21].

### 3.3. Optical properties

$\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) possess different colors.  $\text{EuLn}_2\text{Se}_4$  ( $\text{Ln} = \text{Tb, Dy, Ho, Er, Tm, Lu}$ ) are red.  $\text{EuYb}_2\text{S}_4$  is dark red and  $\text{EuYb}_2\text{Se}_4$  is black. The UV–vis–NIR diffuse reflectance spectra of  $\text{EuLn}_2\text{Q}_4$  are presented in Fig. 6. The band gaps of  $\text{EuTb}_2\text{Se}_4$ ,  $\text{EuDy}_2\text{Se}_4$ ,  $\text{EuHo}_2\text{Se}_4$ ,  $\text{EuEr}_2\text{Se}_4$ ,  $\text{EuTm}_2\text{Se}_4$ ,  $\text{EuYb}_2\text{S}_4$ ,  $\text{EuLu}_2\text{Se}_4$ , and  $\text{EuYb}_2\text{Se}_4$  are found to be 2.0, 1.8, 1.8, 1.7, 1.8, 1.3, 1.7, and 1.6 eV, respectively. These values are consistent with observed colors.  $\text{EuYb}_2\text{Se}_4$  has a much smaller gap than that of other selenides. In addition,  $\text{EuTm}_2\text{S}_4$  and  $\text{EuLu}_2\text{S}_4$ , prepared by the same method, have the same colors and similar band gaps as their selenide analogs [22]. These unusual results conflict with the conventional scheme of the band gap for sulfides and selenides, which is the gap between the valence band of  $3p/4p$  states of S/Se and the conduction band of  $5d$  ( $6s$ ) states of  $\text{Ln}^{3+}$  cations [23,24]. Choices of  $Q$  should play a more

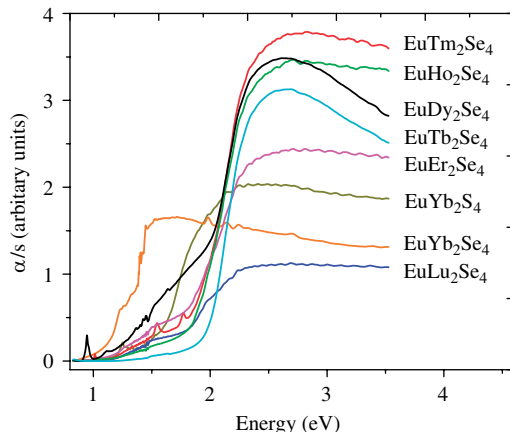


Fig. 6. UV–vis–NIR diffuse reflectance spectra of  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ).

important role than  $\text{Ln}^{3+}$  ions for the band gap values of  $\text{EuLn}_2\text{Q}_4$ . In contrast, the band gap for  $\text{EuYb}_2\text{S}_4$  is larger than the value of  $\text{EuYb}_2\text{Se}_4$  because of the lower energy of  $3p$  states of S. The band gap structures of Yb-based compounds appear to be different from the rest of the compositions. Future work including theoretical calculations for their electronic structures may help to clear up this mystery. The fine structure observed in the spectra for  $\text{EuDy}_2\text{Se}_4$  and  $\text{EuTm}_2\text{Se}_4$  is due to  $f$ – $f$  transitions within the lanthanide ions. There is some tailing below the initial absorption edge for  $\text{EuDy}_2\text{Se}_4$ ,  $\text{EuEr}_2\text{Se}_4$ ,  $\text{EuTm}_2\text{Se}_4$ ,  $\text{EuYb}_2\text{Se}_4$ , and  $\text{EuLu}_2\text{Se}_4$ , indicating the presence of a low-level impurity ( $\text{Sb}_2\text{Se}_3$ ) or indirect band gaps for these materials.

## 4. Conclusions

$\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ ) were prepared using  $\text{Sb}_2\text{Q}_3$  ( $\text{Q} = \text{S, Se}$ ) fluxes and characterized by single-crystal X-ray diffraction, magnetic susceptibility measurements, and UV–vis–NIR diffuse reflectance spectroscopy. This isotopic series adopts the  $\text{CaFe}_2\text{O}_4$ -type three-dimensional channel structure. All of these compounds are antiferromagnets at low temperature due to the magnetic couplings between  $\text{Eu}^{2+}$  ions.  $\text{EuLn}_2\text{Q}_4$  are semiconductors and some of them have atypical band gaps.

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Table 4  
Magnetic parameters for  $\text{EuLn}_2\text{Q}_4$  ( $\text{Ln} = \text{Tb–Lu}$ ;  $\text{Q} = \text{S, Se}$ )

Formula	$P_{\text{cal}} (\mu_B)^a$	$P_{\text{eff}} (\mu_B)^a$	$\theta_p (\text{K})^b$	$R^2$	$T_N (\text{K})$
$\text{EuTb}_2\text{Se}_4$	15.87	16.44(2)	−12.9(6)	0.99984	3.6
$\text{EuDy}_2\text{Se}_4$	17.00	16.315(6)	−3.8(1)	0.99999	3.0
$\text{EuHo}_2\text{Se}_4$	16.96	16.46(2)	−2.7(5)	0.99988	3.2
$\text{EuEr}_2\text{Se}_4$	15.72	15.519(8)	−6.1(2)	0.99998	4.0
$\text{EuTm}_2\text{Se}_4$	13.33	13.35(1)	−7.2(4)	0.99992	4.0
$\text{EuYb}_2\text{S}_4$	10.21	9.563(8)	−9.9(4)	0.99994	4.3
$\text{EuLu}_2\text{Se}_4$	7.94	8.00(1)	−2.4(5)	0.9999	4.4
$\text{EuYb}_2\text{Se}_4$	10.21	8.76(1)	−11.0(5)	0.99989	4.3

<sup>a</sup> $P_{\text{cal}}$  and  $P_{\text{eff}}$ : calculated [19] and experimental effective magnetic moments per formula unit.

<sup>b</sup>Weiss constant ( $\theta_p$ ) and goodness of fit ( $R^2$ ) obtained from high-temperature (100–300 K) data.

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### Appendix A. Supplementary materials

Supplementary data associated with this article can be found in the online version at doi:10.1016/j.jssc.2007.10.028.

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